

PHYSICS END-OF-SEMESTER TEST

The three exercises are independent. One double-sided handwritten A4 form and calculator are allowed. **Any unjustified result will not be taken into account.** Make sure the copy is well presented and legible.

A formula sheet is provided at the end of this document, useful for all three exercises.

1 Modeling a high-voltage line

In simplified terms, a high-voltage line can be seen as a system consisting of a cable (a cylinder) and the earth (a plane), as shown in the figure 1 on the left. The permittivity at any point in the system is fixed at ε_0 .

1.1 Study of the cable without earth

Initially, we will look at the cable alone (without the earth) modelled by a solid cylinder of infinite length, radius a and carrying a uniform surface load $\sigma > 0$.

Question 1-1 : By defining a suitable coordinate system, determine the electric field $\vec{E}$ that it creates throughout space.

Question 1-2 : Deduce the electric potential V at any point in space. It will be assumed to be equal to V_0 at $r = a$.

1.2 Modelling the high-voltage line

To study the high-voltage line, we will use a model consisting of two parallel cables, similar to the previous one, one carrying a surface charge σ and the other $-\sigma$ (Figure 1, right). It is assumed that the distance $2h$ separating the two cables is very large compared with the radius a of the cable. The Cartesian coordinate system is that shown in Figure 1.

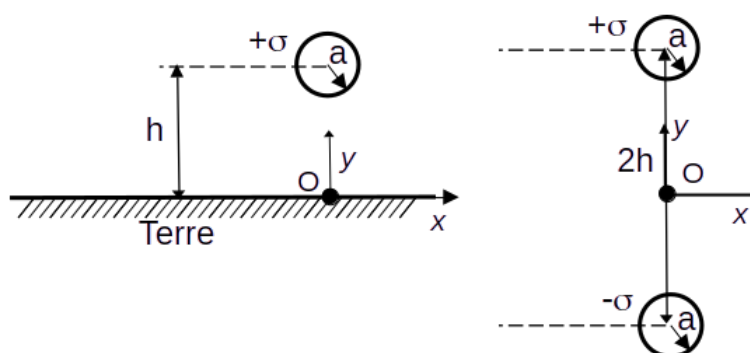


FIGURE 1 : High-voltage line above Earth (left) and associated electrostatic model (right)

Question 1-3 : Study the symmetries and invariances of this new system.

We place ourselves in the plane such that $x = 0$, and with $0 \leq y \leq h$.

Question 1-4 : From the result obtained in question 1.1, express the electric field produced by each cable as a function of the variable y in the portion of plane considered.

Question 1-5 : Deduce the total electric field $\vec{E}_T$ of this system in the considered area.

In relation with electrostatic induction, the earth acquires a charge density σ_T due to the presence of the cable.

Question 1-6 : Calculate σ_T at any point with coordinates $(0, 0, z)$ assuming that the earth is a perfect conductor.

2 Asynchronous motor or induction motor

An asynchronous motor is an AC machine that converts electrical energy into mechanical energy. It consists of a moving circuit, called the **rotor**, which rotates inside a fixed circuit, called the **stator**.

The model shown in Figure 2 consists of a moving **closed rectangular loop or current loop** (the rotor, at the center of the figure), immersed in the magnetic fields $\vec{B}_1(t) = B_0 \cos(\omega_0 t) \vec{u}_x$ and $\vec{B}_2(t) = B_0 \sin(\omega_0 t) \vec{u}_y$ produced respectively by the two stator coils (S_1) and (S_2). The superposition of these two fields creates a **rotating magnetic field** $\vec{B}$, of **constant norm**, but **whose direction rotates** in the (xOy) plane around the axis Oz and forms an angle $\omega_0 t$ with the direction $\vec{u}_x$ at time t (see figure 2).

Throughout the problem, the rectangular loop, of normal $\vec{n}$ contained in the (xOy) plane, rotates around the Oz axis in the same direction of rotation as $\vec{B}(t)$, but **at a constant angular velocity ω different from ω_0** . $\vec{B}(t)$ is uniform over the entire (small) surface S of the loop. The rectangular loop has an **electrical resistance R** and a **self-inductance L** .

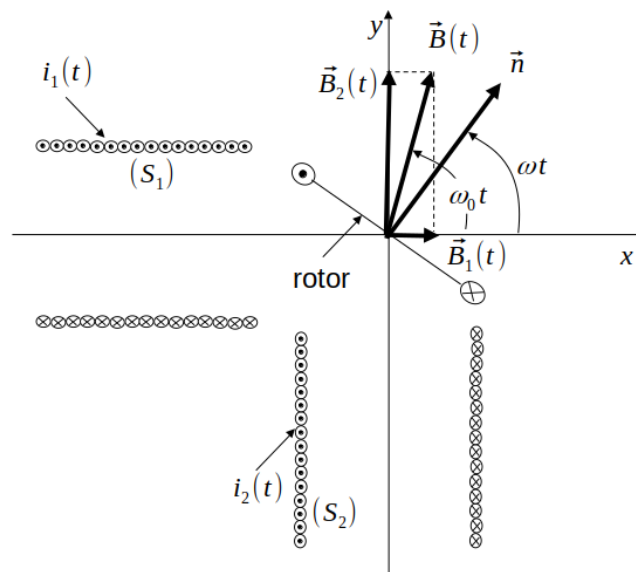


FIGURE 2: asynchronous machine

2.1 Principle of the asynchronous machine

In the entire problem, we're focusing on the steady state, i.e. when the angular velocity ω is constant. In this harmonic regime, **time-varying electrical quantities can be expressed as complex quantities**.

Question 2-1 : Express the flux $\Phi(t)$ of the magnetic field passing through the rectangular loop (or rotor) as a function of B_0 , S , ω_0 and ω .

Question 2-2 : Derive the expression of the induced electromotive force (e.m.f.) $e_{ind}(t)$.

Question 2-3 : Put the expression $e_{ind}(t)$, which is the complex expression of $e_{ind}(t)$, into the form : $e_{ind} = e_m e^{j(\omega_e t - \varphi_e)}$, and deduce the expressions of e_m , φ_e and ω_e .

Question 2-4 : Write the electrical equation that applies to the rectangular loop. Write this equation in the form : $e_{ind} = Z i_{ind}(t)$, with $Z = |Z|e^{j\varphi}$ and deduce the expressions for Z , $|Z|$ and φ as a function of R , L , and ω_e .

Question 2-5 : Deduce the real expression $i_{ind}(t)$ of $i_{ind}(t)$, as a function of e_m , $|Z|$, φ , ω_e and φ_e .

Question 2-6 : Define the magnetic moment $\vec{m}$ of the rectangular loop through which a current i flows. Indicate the relevant quantities on a diagram.

Question 2-7 : Explain qualitatively why the rectangular loop is driven into rotation when the coils (S_1) and (S_2) are energized and produce the rotating magnetic field $\vec{B}(t)$.

Question 2-8 : Express the instantaneous moment $\mathcal{M}_z(t)$ of the magnetic forces acting on the loop of Figure 2 with respect to the axis Oz as a function of e_m , S , B_0 , L , R and ω_e .

Question 2-9 : Express the time average $\langle \mathcal{M}_z \rangle$ of $\mathcal{M}_z(t)$ as a function of the same parameters as for question 2-8.

2.2 Asynchronous motor operating conditions running with no load

In the following questions, you can use Figure 3, which shows the evolution of the normalized mean moment (torque) applied to the rotor as a function of its angular velocity.

Question 2-10 : Which range of angular velocities ω of the rectangular loop corresponds to operation in motor mode of the machine? What can be said about the starting torque ($\omega = 0$)?

Question 2-11 : For $\omega = \omega_0$, what is $\langle \mathcal{M}_z \rangle$? Interpret this value.

Question 2-12 : (BONUS) What happens to the rotational speed of the rotor if the motor is under load with a resisting moment $\langle \mathcal{M}_{z_{res.}} \rangle = -0.2 \langle \mathcal{M}_{z_{max}} \rangle$? Why is it called an **asynchronous** motor?

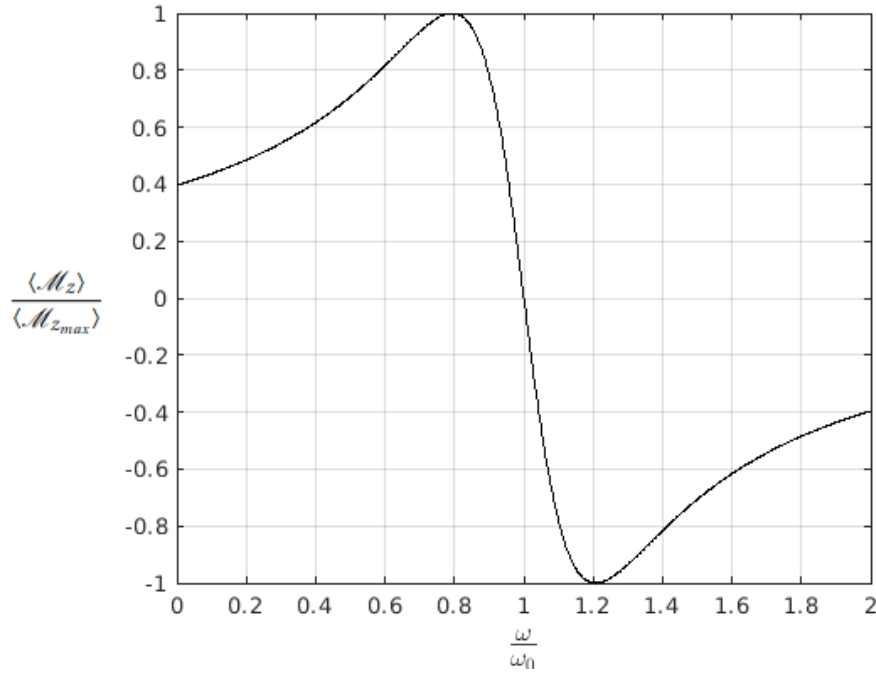


FIGURE 3: Evolution of average moment (torque) as a function of rotor angular velocity

3 Charging cars on inductive roads

3.1 Creation of the magnetic field of the inductor

Let's consider an electric car, equipped with an induction charging system¹: the car drives over a coil buried in the asphalt to recover its autonomy (see figure 4).

All numerical values are gathered at the end of the exercise.

The inductor coil buried in the road consists of N_1 turns through which a current i_1 flows, each describing a rectangle of length D and width W . The stack of N_1 turns spreads over a height H (see figure 4).

First, we'll simplify the coil in order to evaluate the magnetic field it produces. To do this, we choose to **neglect the width W of the coil** (and therefore the inductive effect of portions of the coil perpendicular to the road) and to **reduce it to two planes parallel** to the (xOz) plane, **assumed infinite** in directions parallel to $\vec{u}_x$ and $\vec{u}_z$, and intersecting the (Oy) axis at $y = \frac{W}{2}$ and $y = -\frac{W}{2}$ (see figure 5). Surface currents flow through these planes such that :

$$\vec{k} = +k \vec{u}_x \text{ for } y = -\frac{W}{2} \quad (1)$$

$$\vec{k}' = -k \vec{u}_x \text{ for } y = +\frac{W}{2} \quad (2)$$

with $k > 0$. In the whole space, permeability is μ_0 .

1. This system is still at the experimental stage. See <https://www.incit-ev.eu/> for more information

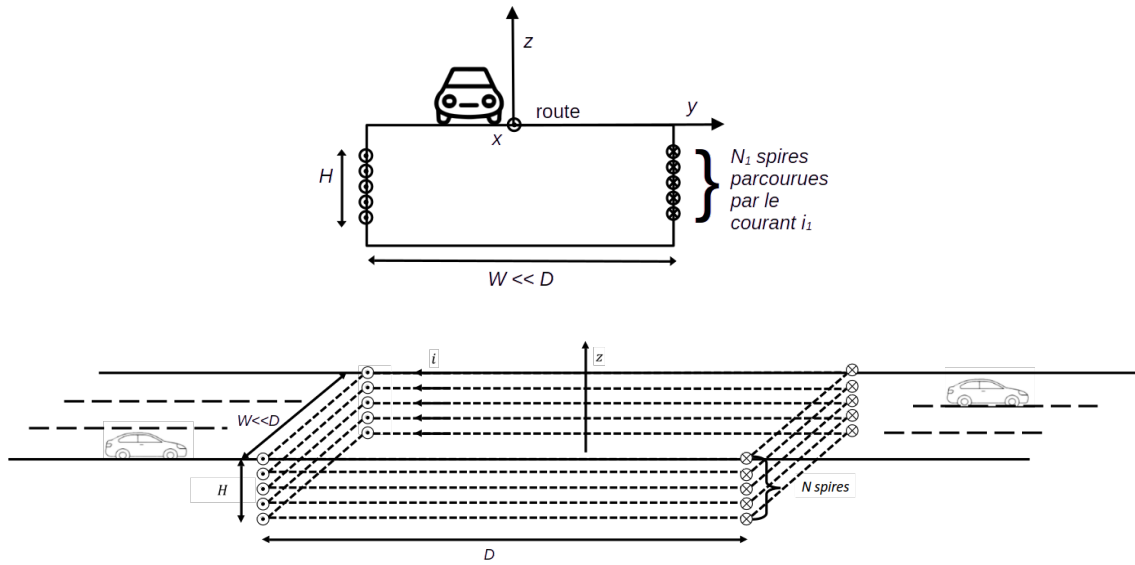


FIGURE 4: Coil structure buried in the road. Parts of the coil along the W width will be ignored

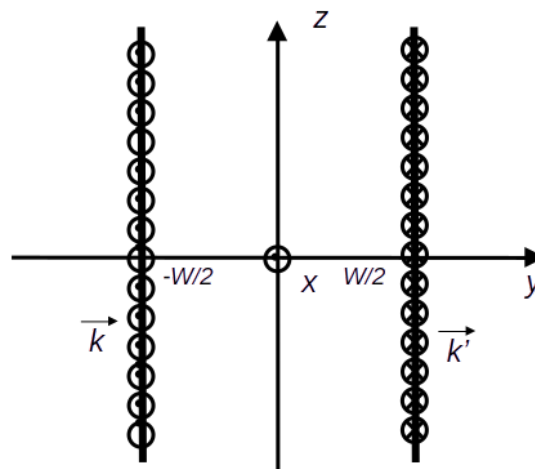


FIGURE 5: Coil model : two parallel planes with a surface current density

Question 3-1 : Determine the relationship between N_1 , i_1 , H and k .

Question 3-2 : From the symmetries and invariances of this global distribution of surface current consisting of the two infinite planes, determine the topology of the magnetic field $\vec{B}_1$ they create.

Question 3-3 : Using symmetry considerations, give a relationship between the magnetic field components of $\vec{B}_1(y = y_0)$ and $\vec{B}_1(y = -y_0)$, $y_0 \in \mathbb{R}^*$.

We now consider only one of the planes, specifically the one passing through $y = -\frac{W}{2}$.

Question 3-4 : By a rigorous and complete study, determine the magnetic field $\vec{B}_0$ created by this single plane in all space.

Question 3-5 : Deduce the expression of the magnetic field $\vec{B}_1$ created in the whole space **by the two planes**. Check that the result agrees with that of questions 3.2 and 3.3.

Question 3-6 : Assuming that the magnetic field of the real (non-infinite) coil is the same as that studied previously, deduce the self-inductance L_1 of the inducting coil as a function of H , W , D , N_1 and μ_0 .

The magnetic field of the coil studied previously is picked up by a flat square coil, **of vertical normal**, placed under the car. It forms the induction receiver of this system and will be referred to as “circuit 2”. This induced coil, with surface area S , comprising N_2 turns, self-inductance L_2 , and resistance R_2 , is connected to an accumulator, modelled by a capacitor of capacitance C_2 . The current i_1 flowing in the inductor coil is now sinusoidal, of angular frequency ω , so that the magnetic field captured by the induced coil is of the form :

$$\vec{B}_1 = B_1 \sqrt{2} \cos(\omega t) \vec{u}_z$$

This magnetic field will be assumed to be **uniform** throughout circuit 2.

The car travels at a constant speed $\vec{v} = v_0 \vec{u}_x$. In the following, we'll assume that all of circuit 2 is immersed in the magnetic field of the inductor.

Question 3-7 : Justify that motion induction doesn't play any role in this situation.

Question 3-8 : Determine the expression of the electromotive force (e.m.f.) $e_2(t)$ induced by the inductor on circuit 2 when the entire surface S of the induced coil is above the inductor coil. Express this voltage in complex notation : $\underline{e}_2(t)$.

Question 3-9 : Draw an electrical diagram of the induced circuit including the accumulator. Give the complex literal expression of the induced current $\underline{i}_2$.

In the rest of the exercise, we set the conditions such that $L_2 C_2 \omega^2 = 1$.

Question 3-10 : What effect does this assumption have on the relationship between the current $\underline{i}_2$ flowing in the induced circuit and the voltage $\underline{e}_2$?

Question 3-11 : The generic expression for the active power P_a developed by a section of circuit through which a current $\underline{i}$ flows and across which a voltage $\underline{u}$ is measured, is :

$$P_a = U_{\text{eff}} I_{\text{eff}} \cos \varphi$$

where U_{eff} and I_{eff} are the RMS amplitudes of voltage $\underline{u}$ and current $\underline{i}$, and φ the phase shift between these two quantities.

Give the numerical value of P_2 , the active power supplied by the induction phenomenon in circuit 2.

Question 3-12 :

Assuming that all the electrical energy accumulated when crossing the recharging zone is recovered and stored in the battery, numerically evaluate this quantity of energy, and the number of additional kilometers of range gained by the car.

Question 3-13 : Suggest improvements to this system (bonus).

Numerical data for this exercise :

- Car velocity v_0 : 110 km/h
- Dimensions of the inductor coil : $D = 1000$ m, $H = 10$ cm, $W = 1$ m
- RMS value of the magnetic field picked up by the receiving coil : 0.24 mT.
- Number of turns of the inductor coil : $N_1 = 1000$

- Number of turns of the induced coil : $N_2 = 60$
- Surface area of the induced coil : $S = 0.1 \text{ m}^2$.
- Frequency of the current in the inductor : 85 kHz .
- Resistance of the induced circuit : $R_2 = 10 \Omega$
- As reminder : $\mu_0 = 4\pi 10^{-7} \text{ H m}^{-1}$
- Energy consumption of an electric car : 20 kWh for 100 km

Formula sheet (Note : not all formulas are useful)

Divergence in cylindrical coordinates :

$$\operatorname{div} \vec{E} = \vec{\nabla} \cdot \vec{E} = \frac{1}{r} \frac{\partial}{\partial r} (r E_r) + \frac{1}{r} \frac{\partial E_\theta}{\partial \theta} + \frac{\partial E_z}{\partial z}$$

Divergence in spherical coordinates :

$$\operatorname{div} \vec{E} = \vec{\nabla} \cdot \vec{E} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 E_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta E_\theta) + \frac{1}{r \sin \theta} \frac{\partial E_\phi}{\partial \phi}$$

Rotational in cylindrical coordinates :

$$\vec{rot} \vec{E} = \vec{\nabla} \wedge \vec{E} = \left(\frac{1}{r} \frac{\partial E_z}{\partial \theta} - \frac{\partial E_\theta}{\partial z} \right) \vec{u}_r + \left(\frac{\partial E_r}{\partial z} - \frac{\partial E_z}{\partial r} \right) \vec{u}_\theta + \left(\frac{1}{r} \frac{\partial}{\partial r} (r E_\theta) - \frac{1}{r} \frac{\partial E_r}{\partial \theta} \right) \vec{u}_z$$

Rotational in spherical coordinates :

$$\begin{aligned} \vec{rot} \vec{E} = \vec{\nabla} \wedge \vec{E} = & \left\{ \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (E_\varphi \sin \theta) - \frac{\partial E_\theta}{\partial \varphi} \right] \right\} \vec{u}_r \\ & + \left\{ \frac{1}{r} \left[\frac{1}{\sin \theta} \frac{\partial E_r}{\partial \varphi} - \frac{\partial}{\partial r} (r E_\varphi) \right] \right\} \vec{u}_\theta \\ & + \left\{ \frac{1}{r} \left[\frac{\partial}{\partial r} (r E_\theta) - \frac{\partial E_r}{\partial \theta} \right] \right\} \vec{u}_\varphi \end{aligned}$$

Trigonometric formulas :

$$\begin{aligned} \cos a \cdot \cos b &= \frac{1}{2} [\cos(a - b) + \cos(a + b)] \\ \sin a \cdot \sin b &= \frac{1}{2} [\cos(a - b) - \cos(a + b)] \\ \cos a \cdot \sin b &= \frac{1}{2} [\sin(a + b) - \sin(a - b)] \end{aligned}$$

Note pour tout le sujet : pas de points si résultat sans unité.

Champ électrique créé par deux fils infinis chargés	Points	Total/ ques- tion
Q1-1. On se place dans un repère cylindrique d'axe z correspondant à l'axe du cylindre Tout plan $(M, \vec{u}_r, \vec{u}_z)$ est plan de symétrie de la distribution de charges, donc de $\vec{E}$ et le champ $\vec{E}(M)$ est contenu dans ce plan, donc $\vec{E} \cdot \vec{u}_\theta = E_\theta = 0$ Tout plan $(M, \vec{u}_r, \vec{u}_\theta)$ est plan de symétrie de la distribution de charges (fils infinis), et $E_z = 0$ Finalement $\vec{E} = E_r(r, \theta, z) \vec{u}_r$ Invariances en θ et z , ce qui aboutit à : $\vec{E} = E_r(r) \vec{u}_r$	0,5 1,5 (-1 par erreur) 1	3
On utilise l'équation de Maxwell-Gauss : $\text{div} \vec{E} = \frac{\rho}{\epsilon_0}$ Ici, $\rho = 0$ car on n'a que des charges surfaciques Avec la divergence en coordonnées cylindriques, il vient : $\frac{1}{r} \frac{\partial}{\partial r}(r E_r) = 0 \Rightarrow r E_r(r > a) = A$ et $r E_r(r < a) = A_0$ avec A et A_0 des constantes Avec $r = 0$ ou $\vec{E}(O) = \vec{0}$ (car xOy , yOz et xOz sont plans de symétrie des charges, donc de $\vec{E}$, d'où $E_z(O) = 0 = E_x(O) = E_y(O)$ respectivement) : $A_0 = 0$ relation de passage en $r = a$: $\Delta(\epsilon \vec{E}_\parallel) = \sigma \vec{u}_r \Rightarrow \epsilon_0 E_r(a^+) - 0 = \epsilon_0 \frac{A}{a} = \sigma$: $\vec{E}(r > a) = \frac{a\sigma}{\epsilon_0 r} \vec{u}_r$	1 1,5 1 1,5	5
Q1-2. A partir de $\vec{E} = -\vec{\nabla} = -\frac{\partial V}{\partial r}$ On détermine V par intégration. Condition aux limites : $V(r \leq a) = V_0$. On trouve $V(r \geq a) = V_0 - \frac{a\sigma}{\epsilon_0} \ln\left(\frac{r}{a}\right)$ par continuité en $r = a$	1 1	2
Q1-3. invariance sur z uniquement (fils infinis) : 0.5. (xOz) plan d'antisymétrie des charges : 0.5. (yOz) plan de symétrie des charges : 0.5. $\forall M, (xMy)$ plan de symétrie des charges et de $\vec{E}$: 0.5 Bonus : $\vec{E} = E_y \vec{u}_y$ dans le plan $x = 0$	2 Bonus : 2	2 (+2 Bo-nus)
Q1-4. avec $\vec{E}_1$ créé par $\sigma_1 = +\sigma$ et $\vec{E}_2$ créé par $\sigma_2 = -\sigma$: $\vec{E}_1(r_1 = h - y) = \frac{a\sigma}{\epsilon_0(h - y)}(-\vec{u}_y)$ et $\vec{E}_2(r_2 = h + y) = \frac{a(-\sigma)}{\epsilon_0(h + y)}(+\vec{u}_y)$	2 × 1	2
Q1-5. Th. de superposition : $\vec{E}(0, y, z) = \vec{E}_1(r_1) + \vec{E}_2(r_2)$ Il vient $\vec{E}(0, y, z) = \frac{a\sigma}{\epsilon_0} \frac{2h}{h^2 - y^2}(-\vec{u}_y)$	0,5 0,5	1
Q1-6. Th. de Coulomb ou bien Relations de passage en M_0 à la surface du conducteur : $\vec{E}(M_0^-) = \vec{0}$ puisque dans conducteur à l'équilibre électrostatique. $\Delta(\vec{E}_\parallel) = \vec{0} \Rightarrow \vec{E}_\parallel(M_0^+) = \vec{0}$ $\Delta(\epsilon \vec{E}_\perp) = \sigma_T \vec{n} \Rightarrow \vec{E}(M_0^+) = E_\parallel(M_0^+) + \vec{E}_\perp(M_0^+) = \vec{E}_\perp(M_0^+) = \frac{\sigma_T}{\epsilon_0} \vec{n}$ $(0, 0, z)$ est à la surface de la Terre considérée comme un conducteur, de normale $\vec{n} = \vec{u}_y$ $\frac{\sigma_T}{\epsilon_0} \vec{u}_y = \vec{E}(0, 0, z) = -\frac{2\sigma}{\epsilon_0} \frac{a}{h} \vec{u}_y \Rightarrow \sigma_T = -\frac{2a\sigma}{h}$	1 2 1	4
Total pour cette partie :		19 + 2 bonus

Moteur asynchrone	Points	Total/ ques- tion
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Q2-1. $\Phi(t) = \iint_{\text{spire}} \vec{B} \cdot d\vec{S}$, comme le champ est considéré comme uniforme sur la spire cela donne $B_0 S \cos((\vec{B}, \vec{n})) \Rightarrow \Phi(t) = B_0 S \cos((\omega_0 - \omega)t)$ bonus : ici on n'utilise $\vec{B} = \vec{B}^{app}$ car e_{ind}^p est pris en compte avec la donnée de L de la spire	2	
Q2-2. $e_{ind} = -\frac{d\Phi(t)}{dt} = (\omega_0 - \omega)B_0 S \sin((\omega_0 - \omega)t)$.	1	1
Q2-3. $e_m = (\omega_0 - \omega)B_0 S$, $\varphi_e = \frac{\pi}{2}$ et $\omega_e = \omega_0 - \omega$	0,5+1+0,5	2
Q2-4. $e_{ind}(t) = Ri(t) + L \frac{di(t)}{dt}$ En régime harmonique : $e_{ind}(t) = (R + jL\omega_e)i_{ind}(t)$ d'où $Z = R + jL\omega_e$, $ Z = \sqrt{R^2 + (L\omega_e)^2}$ et $\tan \varphi = \frac{L\omega_e}{R}$	1 1	2
Q2-5. $i_{ind} = \frac{e_{ind}}{Z} = \frac{e_m}{ Z } e^{j(\omega_e t - \frac{\pi}{2} - \varphi)} \Rightarrow i_{ind}(t) = \text{Re}(i_{ind}(t)) = \frac{e_m}{ Z } \cos(\omega_e t - \frac{\pi}{2} - \varphi)$ $\Rightarrow i_{ind}(t) = \frac{e_m}{ Z } \sin(\omega_e t - \varphi)$ remarque : le calcul se fait aussi à partir des représentations algébriques des complexes : $i_{ind} = \frac{e_{ind}}{Z} = \frac{e_m}{R^2 + L^2\omega_e^2} (R - jL\omega_e)(-j) e^{j(\omega_e t)}$ d'où $i_{ind} = \text{Re}(i_{ind}) = \frac{e_m}{R^2 + L^2\omega_e^2} (R \sin(\omega_e t) - L\omega_e \cos(\omega_e t))$	1,5	1,5
Q2-6. $\vec{m} = iS\vec{n}$. Le schéma doit montrer le sens positif choisi pour le courant et la normale orientée qui en résulte.	1 (-0,5 par oubli)	1
Q2-7. La variation de flux magnétique à travers la spire, due à la rotation du champ magnétique engendre une f.é.m. induite (induction statique), elle même à l'origine d'un <u>courant induit</u> . La spire parcourue par un courant et plongée dans le champ magnétique uniforme est soumise à des <u>forces de Laplace</u> dont l'action la met en mouvement de rotation. Elle est soumise au <u>couple de moment résultant</u> $\vec{\mathcal{M}}(t) = \vec{m}(t) \wedge \vec{B}(t)$. On peut aussi invoquer la règle du flux maximum et dire que la boucle de courant abandonnée aux forces de Laplace évolue de manière à maximiser le flux. Elle cherche donc à aligner sa normale avec la direction du champ magnétique. Donc elle tourne et cherche à la rattraper.	2 (-0,5 par oubli)	2
Q2-8. $\vec{\mathcal{M}}(t) = i_{ind} S \vec{n} \wedge \vec{B} = i_{ind} S B_0 \sin(\omega_e t) \vec{u}_z$ $\Rightarrow \mathcal{M}_z(t) = \frac{e_m S B_0}{ Z } \sin(\omega_e t - \varphi) \sin(\omega_e t)$ (ou encore $= \frac{e_m S B_0}{2 Z } [\cos \varphi - \cos(2\omega_e t - \varphi)]$) (Remarque : Comme $\cos \varphi = \frac{R}{ Z }$ et $\sin \varphi = \frac{L\omega_e}{ Z }$, après développement on obtient par ex. $\mathcal{M}_z = \frac{e_m S B_0}{R^2 + L^2\omega_e^2} [R \sin(\omega_e t) - L\omega_e \cos(\omega_e t)] \sin(\omega_e t)$ mais l'expression de la ligne précédente suffit) bonus : ici on ne prend pas en compte $\vec{B}_p$ dans $\vec{\mathcal{M}}$ car contribue uniquement au torseurs des forces intérieures, or $\vec{\mathcal{M}}_{int} = \vec{0}$	2 (ou toute expression équivalente) bonus 0,5	2
Q2-9. $\langle \cos(2\omega_e t - \varphi) \rangle = 0$, ou bien $\int_0^{T_e} \cos(\omega_e t) \sin(\omega_e t) dt = 0$ et $\frac{1}{T_e} \int_0^{T_e} \sin^2(\omega_e t) dt = \frac{1}{2}$, donc $\langle \mathcal{M}_z \rangle = \frac{e_m S B_0}{2 Z } \cos \varphi = \frac{\omega_e (S B_0)^2 \cos \varphi}{2} = \frac{R \omega_e (S B_0)^2}{2(R^2 + L^2\omega_e^2)}$	1,5	1,5
Q2-10. Le couple est moteur pour $\langle \mathcal{M}_z \rangle \geq 0 \Leftrightarrow 0 \leq \omega \leq \omega_0$, le rotor est toujours un peu en retard sur le champ magnétique. Au démarrage, le couple est bien moteur et le moteur peut démarrer, sauf si un couple résistant trop important est présent quelque part (supérieur à $0.6 \mathcal{M}_{max}$)	1 bonus 0,5 0,5 bonus 0,5	1,5
Q2-11. Le couple est nul : frontière entre le fonctionnement moteur et génératrice (non exigible : le rotor est accroché sur le champ magnétique.)	1	1

Q2-12. BONUS : Si on a un couple résistant $\langle \mathcal{M}_{z_{rés.}} \rangle = -0.2 \langle \mathcal{M}_{z_{max}} \rangle$, le rotor se stabilisera à une vitesse telle que le couple total est nul, soit $\langle \mathcal{M}_z \rangle + \langle \mathcal{M}_{z_{rés.}} \rangle = 0 \Rightarrow \frac{\langle \mathcal{M}_z \rangle}{\langle \mathcal{M}_{z_{max}} \rangle} = 0.2$. On voit sur la figure que le point de fonctionnement va remonter sur la partie verticale de la courbe jusqu'à une valeur du moment moyen normalisé de 0.2 et que par conséquent la vitesse angulaire ω va prendre une valeur légèrement inférieure à ω_0 . Le rotor tourne donc légèrement moins vite que le champ magnétique, d'où le nom de moteur asynchrone. Surbonus : On remarque qu'en raison de la vitesse relative du rotor (l'induit) par rapport au champ magnétique, on est en présence d'un phénomène d'induction mutuelle.	bonus 1	
Total de cette partie :		17.5 + bonus 3.5

Exercice 3 : Recharge des voitures électrique sur des portions de route inductrices (27,5/63 + bonus 1)	Points	Total/ques-tion
Q3-1. $N_1 i_1 = \iint \vec{j} \cdot \vec{u}_x dS = \int_0^H \vec{k} \cdot \vec{u}_x dz \Leftrightarrow N_1 i_1 = kH$	1	1
Q3-2. dans le modèle nappes de courant parallèles à (xOz) infinies : invariances par toutes translations colinéaires à $\vec{u}_x$ et $\vec{u}_z$: $\vec{B}_1(y)$ (xMy) plan de symétrie des densités de courants $\vec{k}$ et $\vec{k}'$, donc d'antisymétrie de $\vec{B}_1$ et $\vec{B}_1(M) = B_{1z}(y)\vec{u}_z$	0,5 1 (être exigeant)	1,5
Q3-3. (xOz) plan d'antisymétrie des densités de courants, donc de symétrie de $\vec{B}_1$ et $B_{1z}(-y) = B_{1z}(y)$ (les composantes $\parallel$ de $\vec{B}_1$ sont des fonctions paires)	1	1
Q3-4. On a toujours $\vec{B}_0(y)$ par invariances, (xMy) plan d'antisymétrie de $\vec{B}_0$ et $\vec{B}_0(M) = B_{0z}(y)\vec{u}_z$: dès lors $\vec{\text{rot}} \left(\frac{\vec{B}}{\mu} \right) = \frac{1}{\mu_0} \frac{\partial B_{0z}}{\partial y} \vec{u}_x = \frac{1}{\mu_0} \frac{dB_{0z}}{dy} \vec{u}_x \stackrel{MA}{=} \vec{0}$ (pas de courant volumique), d'où $B_{0z}(y > -\frac{W}{2}) = \text{cte1} = B_0^+$ et $B_{0z}(y < -\frac{W}{2}) = \text{cte2} = B_0^-$ par ailleurs $y = -\frac{W}{2}$ plan de symétrie de $\vec{k}$, donc d'antisymétrie de $\vec{B}_0$, d'où $B_0^- = -B_0^+$ (eq1- B_0) Relation de passage sur $\Delta \vec{B}_{\parallel}$ en $y = -\frac{W}{2}$ conduit à $\frac{B_{0z}(y = -\frac{W}{2}^+)}{\mu_0} \vec{u}_z - \frac{B_{0z}(y = -\frac{W}{2}^-)}{\mu_0} \vec{u}_z = k \vec{u}_x \wedge \vec{u}_y = k \vec{u}_z \Leftrightarrow B_0^+ - B_0^- \stackrel{\text{eq1-}B_0}{=} 2B_0^+ = \mu_0 k$	0,5 1,5 0,5 1 0,5 expression générale correcte 1	5
Q3-5. $\vec{B}'_0$ créé par $\vec{k}'$ se déduit de $\vec{B}_0$ en remplaçant k par $k'_x = -k$: $B'_{0z}(y > \frac{W}{2}) = B_0^+ + \frac{\mu_0 k'_x}{2} = -\frac{\mu_0 k}{2} = -B'_{0z}(y < \frac{W}{2})$ on utilise le Th. de superposition en distinguant les zones de l'espace $y < -\frac{W}{2}$, $-\frac{W}{2} < y < \frac{W}{2}$ et $\frac{W}{2} < y$, on trouve : $B_{1,z}(y < -\frac{W}{2}) = B_0^- + B_0'^- = -\frac{\mu_0 k}{2} + \frac{\mu_0 k}{2} = 0$, $B_{1,z}(-\frac{W}{2} < y < \frac{W}{2}) = B_0^+ + B_0'^- = \frac{\mu_0 k}{2} + \frac{\mu_0 k}{2} = \mu_0 k$, et $B_{1,z}(y > \frac{W}{2}) = B_0^+ + B_0'^+ = \frac{\mu_0 k}{2} - \frac{\mu_0 k}{2} = 0$: on retrouve bien que $\vec{B}_1 = B_{1z}(y)\vec{u}_z$ et $B_{1z}(y)$ est paire	1 0,5 1,5 0,5 tout ou rien	3,5
Q3-6. en utilisant $\vec{n} = +\vec{u}_z$ imposé par la définition du sens de i_1	0,5 précision sens $\vec{n}$	

à travers 1 spire : $\varphi_1 = \iint \vec{B}_1 \cdot \vec{u}_z dS = \mu_0 k DW$, et $L_1 = \frac{N_1 \varphi_1}{i_1} = \frac{\mu_0 N_1 k W D}{k \frac{H}{N_1}} = \frac{\mu_0 N_1^2 W D}{H}$	1,5 (-1 $L_1 < 0$ sans comm.)	2
Q3-7. $\vec{E}_m^{mot} = \vec{v} \wedge \vec{B}_1 = -v_0 B_1 \vec{u}_y$ uniforme : dès lors $e_{ind}^{mot} = \oint \vec{E}_m^{mot} \cdot d\vec{\ell} = -v_0 B_1 \vec{u}_y \cdot \oint d\vec{OM} = -v_0 B_1 \vec{u}_y \cdot [\vec{OM}]_A^{B=A} = 0$	1 1	2
Q3-8. avec e_2 calculée dans le sens défini par $\vec{n}_2 = +\vec{u}_z$: $e_2 = -\frac{d(N_2 \varphi_2)}{dt}$, $\varphi_2 = \iint \vec{B}_1 \cdot \vec{u}_z dS = B_{1,z}(t)S$ ($\vec{B}_1$ uniforme). et finalement $e_2 = -j\omega N_2 B_1 S \sqrt{2} e^{j\omega t}$	0,5 précision sens de e_2 calculé 1,5 0,5	2,5
Q3-9 circuit 1 maille avec i_2 de sens e_2 : $i_2 = \frac{e_2}{R_2 + jL_2\omega - \frac{j}{C_2\omega}} = \frac{e_2}{R_2 + j\frac{L_2 C_2 \omega^2 - 1}{C_2 \omega}}$ (éq. circuit2)	0,5 définition sens de i_2 1 (-0,5/ erreur signe)	1,5
Q3-10. (éq. circuit2) devient $i_2 = \frac{e_2}{R_2}$ et i_2 et e_2 sont en phase.	1	1
Q3-11. en phase donc $P_{ind} = i_{2,eff} e_{2,eff} \times 1 = \frac{i_{2,max}}{\sqrt{2}} \frac{e_{2,max}}{\sqrt{2}}$. $i_{2,max} = \frac{e_{2,max}}{R_2}$ et $e_{2,max} = \omega N_2 B_1 S \sqrt{2}$: $P_{ind} = \frac{\omega^2 N_2^2 B_1^2 S^2}{R_2}$. AN : $P_{ind} \simeq 59 \text{ kW}$	1 1 1,5 (0 sans unité)	3,5
Q3-12. en notant d la distance parcourue pour consommer toute l'énergie accumulée, $\eta = \frac{20 \text{ kWh}}{100 \text{ km}}$ la consommation sur 100km, et $\tau = \frac{D}{v_0}$, le temps de parcours sur la portion de route inductrice (en négligeant les durée d'entrée et de sortie du circuit 2 de cette zone, donc en considérant que toute la surface S du circuit 2 est entièrement dans cette zone pendant toute la durée τ) énergie accumulée $P_{ind} \tau = \eta d \Leftrightarrow d = \frac{P_{ind} D}{v \eta}$. AN : Energie accumulée : 1930 kJ (On rappelle que : 1 Wh = 3600 J) $d = \frac{59,1 \cdot 10^3 \text{ W} \times 1 \text{ km}}{110 \text{ km h}^{-1} \times 20 \cdot 10^3 \text{ Wh (100 km)}^{-1}} \simeq 2.7 \text{ km}$	1 (pour toute expression littérale cohérente, être exigeant sur la rédaction) 0.5 1 (-1 sans unité)	2.5
Q3-13. bonus : améliorations possibles Le mieux est d'augmenter si possible toutes les grandeurs qui figurent au carré dans l'expression de la puissance. On peut imaginer augmenter le nombre de spires induites (sous la voiture), mais on augmente le poids et la résistance du circuit. On peut aussi augmenter la fréquence, mais il faut ensuite ajuster les impédances du circuit induit. On peut augmenter la valeur du champ magnétique, mais il faut faire passer un courant important dans l'inducteur...	bonus 2	
Total pour cette partie :		27 + 2 bonus
Total général :		63.5 points + 7.5 bonus